

An EOS implementation for astrophysical simulations

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Nuclear EOS

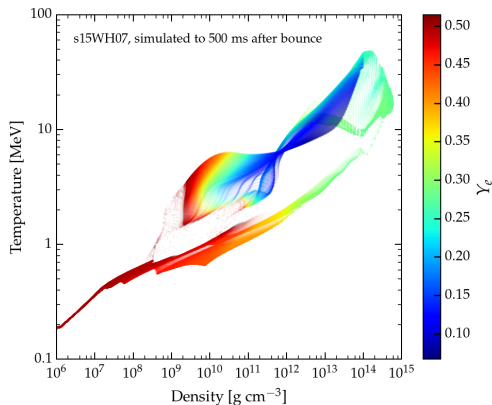
Equation of state (EOS):

$$\Rightarrow \varepsilon(n, y, T), P(n, y, T),$$

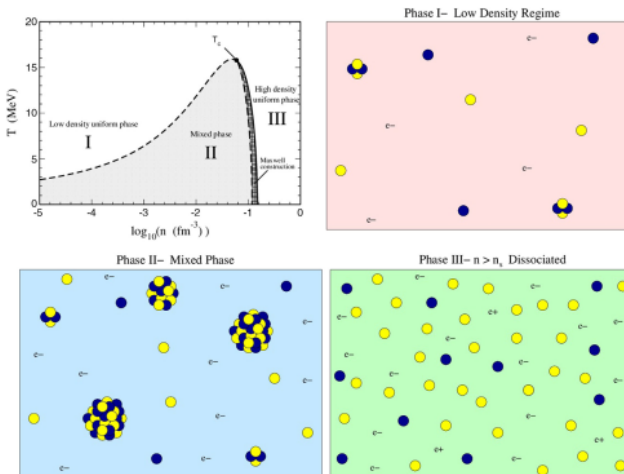
$$s(n, y, T), \dots$$

Astrophysical relevance

- Core-collapse supernovae;
- NS structure and evolution;
- Merger of compact stars;
- r -process nucleosynthesis;
- ...



Nuclear EOS



Nuclear EOS

Long-standing problem in nuclear physics.

Combines efforts from:

- heavy ion collision experiments;
- nuclear reaction experiments;
- computer simulations of astrophysical phenomena;
- computer simulations of dense matter;
- theoretical many-body calculations;
- ...

Nuclear EOS

Nuclear forces are complicated \Rightarrow combine different approaches!

Hot-dense EOS available:

- 1 Lattimer and Swesty
- 2 H. Shen *et al.*
- 3 G. Shen *et al.*
- 4 Hempel *et al.*
- 5 Steiner *et al.*
- 6 Banik *et al.*
- 7 ...

Classification:

- 1 Relativistic vs
Non-relativistic
- 2 Realistic potentials vs
Effective potentials
- 3 SNA vs NSE vs reaction
networks
- 4 Muons? Hyperons?
Quarks?

Nuclear EOS

Goals:

- 1 Write code to construct EOS tables for astrophysical simulations.
- 2 Easy to update EOS as new nuclear matter constraints become available.
- 3 Make the code open-source. (soon)

The Lattimer & Swesty EOS

The Lattimer & Swesty EOS [Nucl. Phys. A 535, 331 (1991)]

Most used EOS for simulations of CCSNe and NS mergers.

Non-relativistic compressible liquid-drop description of nuclei.

Contains

- 1 Nucleons;
- 2 alpha particles;
- 3 electrons and positrons;
- 4 photons.

Nucleons may cluster to form nuclei.

LS EOS use the single nucleus approximation (SNA).

The Lattimer & Swesty EOS

Free energy $F(n, y, T) = F_o + F_h + F_\alpha + F_e + F_\gamma$

- $F_o \equiv$ nucleons outside heavy nuclei (nucleon gas)
- $F_h \equiv$ nucleons clustered into heavy nuclei

In this work, F depends on seven variables:

- u : volume fraction occupied by heavy nuclei
- r : generalized size of heavy nuclei
- n_{ni} : neutron density inside heavy nuclei
- n_{pi} : proton density inside heavy nuclei
- n_{no} : neutron density outside heavy nuclei
- n_{po} : proton density outside heavy nuclei
- n_α : alpha particle density

The Lattimer & Swesty EOS

Heavy nuclei free energy:

$$F_h = F_i + F_S + F_C + F_T$$

- $F_i \equiv$ nucleons inside heavy nuclei
- $F_S \equiv$ surface free energy
- $F_C \equiv$ coulomb free energy
- $F_T \equiv$ translational free energy

The EOS of each component:

- Nucleons \Rightarrow local phenomenological Skyrme-type effective interaction.
- Alpha particles \Rightarrow hard spheres.
- Electrons, positrons and photons \Rightarrow background gas.

The Lattimer & Swesty EOS

Energy density of bulk nuclear matter with Skyrme-type interactions

$$E_B(n, y, T) = \frac{\hbar^2}{2m_n^*} \tau_n + \frac{\hbar^2}{2m_p^*} \tau_p + (a + 4by(1-y)) n^2 \\ + \sum_i (c_i + 4d_i y(1-y)) n^{1+\delta_i} - yn(m_n - m_p).$$

Nucleon effective mass m_t^*

$$\frac{\hbar^2}{2m_t^*} = \frac{\hbar^2}{2m_t} + \alpha_1 n_t + \alpha_2 n_{-t}.$$

where $t = n \Rightarrow -t = p$ and vice versa, $n_n = (1-y)n$ and $n_p = yn$.

Nucleon kinetic energy density τ_t

$$\tau_t = \frac{1}{2\pi^2} \left(\frac{2m_t^* T}{\hbar^2} \right)^{\frac{5}{2}} F_{3/2}(\eta_t(n, y)),$$

The Lattimer & Swesty EOS

$$a = \frac{t_0}{4}(1 - x_0),$$

$$b = \frac{t_0}{8}(2x_0 + 1),$$

$$c_i = \frac{t_{3i}}{24}(1 - x_{3i}),$$

$$d_i = \frac{t_{3i}}{48}(2x_{3i} + 1),$$

$$\delta_i = \sigma_i + 1,$$

$$\alpha_1 = \frac{1}{8} [t_1(1 - x_1) + 3t_2(1 + x_2)],$$

$$\alpha_2 = \frac{1}{8} [t_1(2 + x_1) + t_2(2 + x_2)].$$

Nuclear surface and Coulomb free energies

$$F_S = \frac{3s(u)}{r} \sigma(y_i, T) \quad \text{and} \quad F_C = \frac{4\pi\alpha}{5} (y_i n_i r)^2 c(u).$$

- $s(u)$: surface shape function
- $c(u)$: Coulomb shape function
- r : generalized nuclear size
- $\sigma(y_i, T)$: surface tension per unit area

Nuclear virial Theorem: $F_S = 2F_C$

$$r = \frac{9\sigma}{2\beta} \left[\frac{s(u)}{c(u)} \right]^{1/3} \quad \text{where} \quad \beta = 9 \left[\frac{\pi\alpha}{15} \right]^{1/3} (y_i n_i \sigma)^{2/3}$$

$$F_S + F_C = \beta [c(u)s(u)^2]^{1/3} \equiv \beta \mathcal{D}(u).$$

Nuclear surface and Coulomb free energies

Lattimer & Swesty interpolate $\mathcal{D}(u)$ with

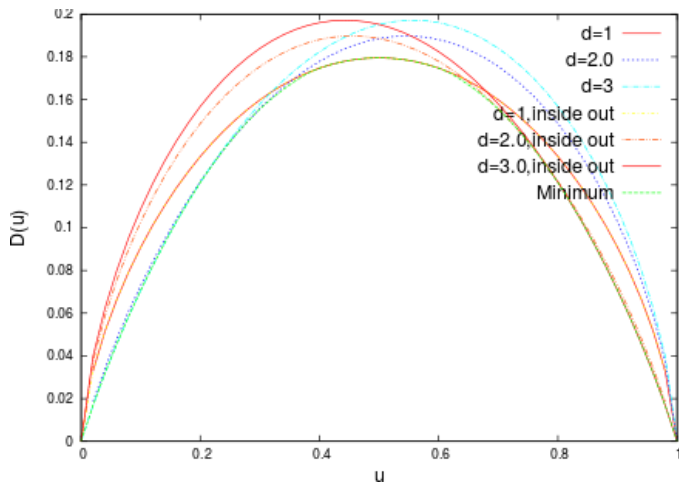
$$\mathcal{D}(u) = u(1-u) \frac{(1-u)D(u)^{1/3} + uD(1-u)^{1/3}}{u^2 + (1-u)^2 + 0.6u^2(1-u)^2}$$

where

$$D(u) = 1 - \frac{3}{2}u^{1/3} + \frac{1}{2}u$$

- as $u \rightarrow 0$ reproduces free energy of spherical nuclei
- as $u \rightarrow 1$ reproduces free energy of “bubble nuclei”
- intermediate u : reproduces free energy of pasta phases:
 - cylinders;
 - slabs;
 - cylindrical holes.

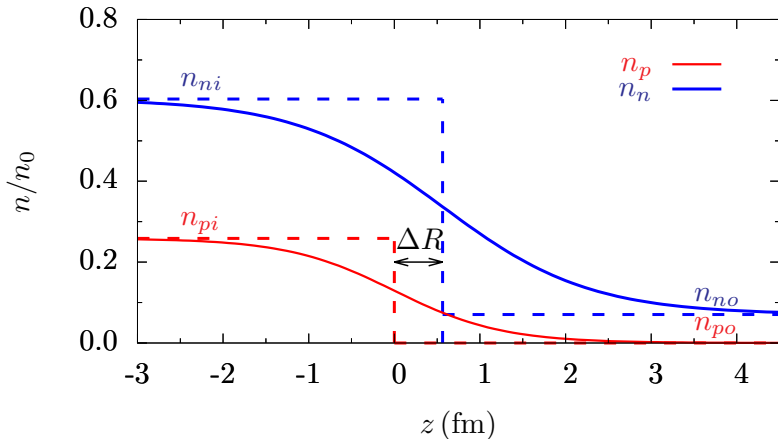
Nuclear surface and Coulomb free energies



Lim (2012)

Nuclear surface and Coulomb free energies

Surface Tension $\sigma(y, T)$.



Nuclear surface and Coulomb free energies

Set temperature T and proton fraction y_i .

Solve equilibrium equations:

$$P_i = P_o, \quad \mu_{ni} = \mu_{no}, \quad \text{and} \quad \mu_{pi} = \mu_{po}, \quad \text{and} \quad y_i = \frac{n_{pi}}{n_{ni} + n_{pi}}$$

Set density to

$$n_t(z) = n_{to} + \frac{n_{ti} - n_{to}}{1 + \exp((z - z_t)/a_t)}$$

$t = n, p$.

Nuclear surface and Coulomb free energies

Find z_t and a_t that minimizes

$$\sigma(y_i, T) = \int_{-\infty}^{+\infty} \left[F_B(z) + E_S(z) + P_o - \mu_{no} n_n(z) - \mu_{po} n_p(z) \right] dz.$$

where

$$E_S(z) = \frac{1}{2} \left[q_{nn} (\nabla n_n)^2 + q_{np} \nabla n_n \cdot \nabla n_p + q_{pn} \nabla n_p \cdot \nabla n_n + q_{pp} (\nabla n_p)^2 \right]$$

and

$$q_{nn} = q_{pp} = \frac{3}{16} [t_1(1 - x_1) - t_2(1 + x_2)],$$

$$q_{np} = q_{pn} = \frac{1}{16} [3t_1(2 + x_1) - t_2(2 + x_2)].$$

Solving the EOS

Minimize $F(n, y, T)$ w.r.t. $u, r, n_{ni}, n_{pi}, n_{no}, n_{po}$, and n_α .

$$A_1 = P_i - B_1 - P_o - P_\alpha = 0,$$

$$A_2 = \mu_{ni} - B_2 - \mu_{no} = 0,$$

$$A_3 = \mu_{pi} - B_3 - \mu_{po} = 0.$$

where

$$B_1 = \frac{\partial \mathcal{F}}{\partial u} - \frac{n_i}{u} \frac{\partial \mathcal{F}}{\partial n_i},$$

$$B_2 = \frac{1}{u} \left[\frac{y_i}{n_i} \frac{\partial \mathcal{F}}{\partial y_i} - \frac{\partial \mathcal{F}}{\partial n_i} \right],$$

$$B_3 = -\frac{1}{u} \left[\frac{1 - y_i}{n_i} \frac{\partial \mathcal{F}}{\partial y_i} + \frac{\partial \mathcal{F}}{\partial n_i} \right],$$

with $\mathcal{F} = F_S + F_C + F_T$.

Solving the EOS

Constraints

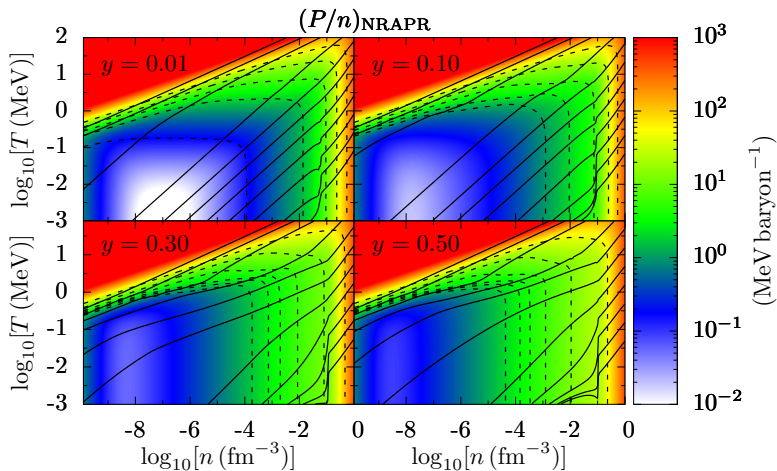
$$n = un_i + (1 - u)[4n_\alpha + n_o(1 - n_\alpha v_\alpha)],$$

$$ny = un_i y_i + (1 - u)[2n_\alpha + n_o y_o(1 - n_\alpha v_\alpha)].$$

$$\mu_\alpha = 2(\mu_{no} + \mu_{po}) + B_\alpha - P_o v_\alpha,$$

$$r = \frac{9\sigma}{2\beta} \left[\frac{s(u)}{c(u)} \right]^{1/3},$$

Solving the EOS



Solving the EOS

Nuclear Statistical Equilibrium

Consider ensemble of nuclei at low densities.

Given an ensemble of nuclei i , solve for μ_n and μ_p such that

$$\begin{aligned}\mu_i &= m_i + E_{c,i} + T \log \left[\frac{n_i}{g_i} \left(\frac{2\pi}{m_i T} \right)^{3/2} \right], \\ &= Z_i \mu_p + (A_i - Z_i) \mu_n\end{aligned}$$

that minimizes the free energy of the system.

Nuclear Statistical Equilibrium

L&S EOS is obtained in the single nucleus approximation (SNA). Properties in the SNA can differ significantly from observed properties of nuclei.

- No shell closure;
- No pairing;
- liquid drop model neglects many-body effects;
- ...

Conversely,

- NSE breaks down close to nuclear saturation density
 $n_0 \simeq 0.16 \text{ fm}^{-3}$;
- Needs very large and very neutron rich nuclei at low y and/or high $n \sim n_0$;
- No nuclear inversion (pasta phase).

Nuclear Statistical Equilibrium

Use ad-hoc procedure to mix NSE and SNA free energies:

$$F_{\text{MIX}} = \chi(n)F_{\text{SNA}} + [1 - \chi(n)]F_{\text{NSE}}.$$

Chose $a(n)$ such that:

$$\chi(n) \rightarrow 0, \text{ if } n \ll n_0$$

$$\chi(n) \rightarrow 1, \text{ if } n \lesssim n_0/10$$

Corrections to thermodynamic quantities, *e.g.*

$$P_{\text{MIX}} = n^2 \left. \frac{\partial(F_{\text{MIX}}/n)}{\partial n} \right|_{T,y} = \chi(n)P_{\text{SNA}} + [1 - \chi(n)]P_{\text{NSE}} + n^2 \frac{\partial\chi(n)}{\partial n} (F_{\text{SNA}} - F_{\text{NSE}}).$$

EOS is self consistent!

High density extension

- Skyrme parametrizations only constrained up to $n \lesssim 3n_0$.
- NS maximum mass depend on EOS at $n \sim 10n_0$.
- Most Skyrme EOS unable to reproduce NS maximum mass, $M_{\max} \sim 2M_{\odot}$.

Add extra c_i , d_i , and δ_i terms to Skyrme interactions:

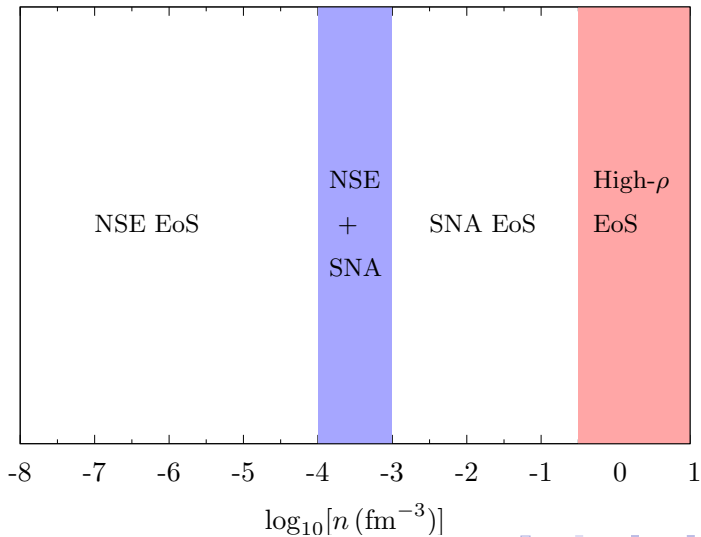
$$\begin{aligned} \varepsilon_B(n, y, T) = & \frac{\hbar^2}{2m_n^*} \tau_n + \frac{\hbar^2}{2m_p^*} \tau_p + (a + 4by(1-y)) n^2 \\ & + \sum_i (c_i + 4d_i y(1-y)) n^{1+\delta_i} - yn(m_n - m_p). \end{aligned}$$

Extra terms should:

- barely affect EOS for $n \lesssim 3n_0$;
- increase $M_{\max} \sim 2M_{\odot}$.

Problem: may imply in $(c_s/c) \gtrsim 1$ for $n \sim (6-10)n_0$.

Final EOS



Skyrme parametrizations

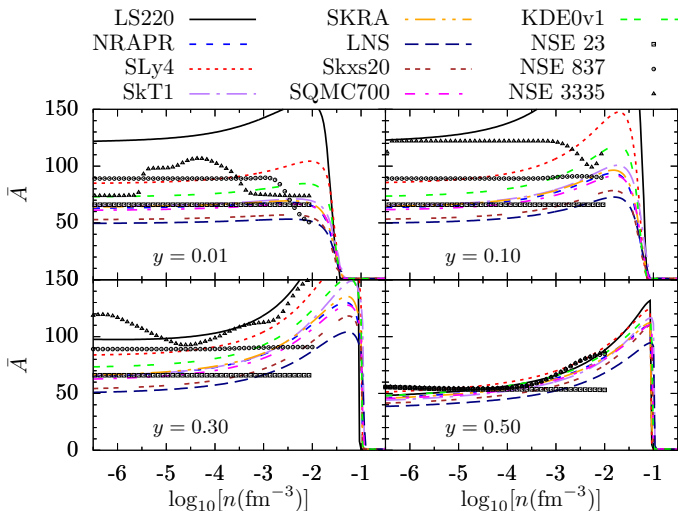
Dutra *et al.* [Phys. Rev. C 85 035201 (2012)]

- Analyzed over 240 Skyrme parametrizations available in the literature.
- Only 11 fulfill all well established nuclear physics constraints!
- Not all 11 reproduce $M_{\max} \sim 2M_{\odot}$!

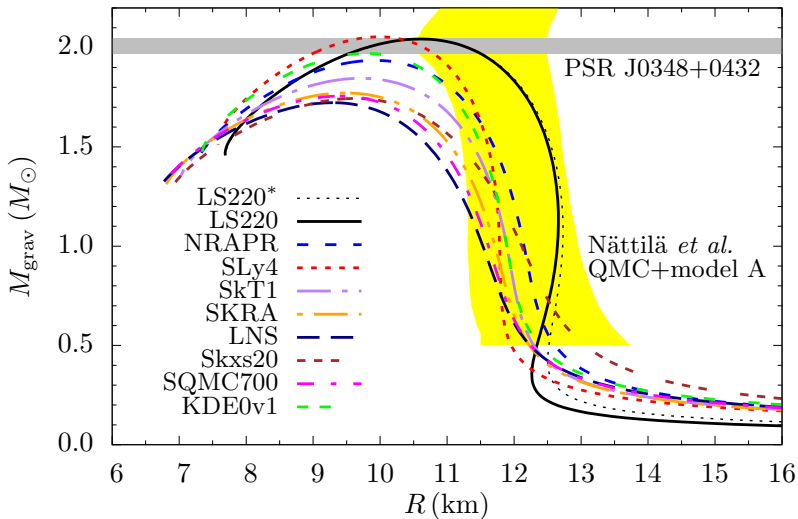
We produced hot-dense EOS tables for a few of these parametrizations.

Skyrme parametrizations

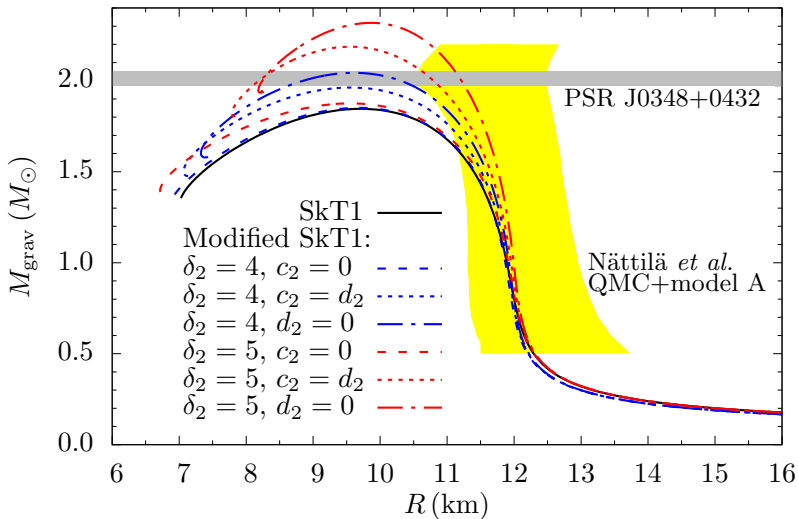
Average nuclear size \bar{A} along $s = 1 k_B \text{baryon}^{-1}$



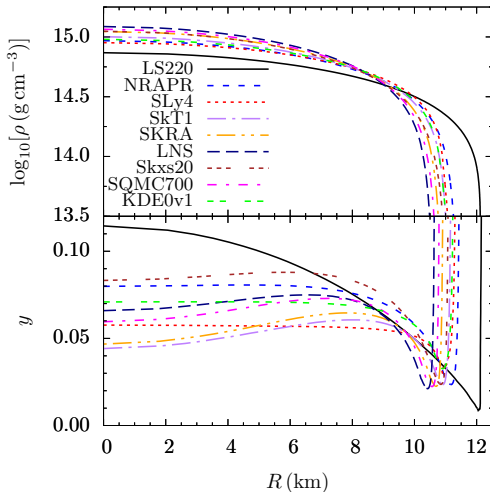
NS mass-radius relationship



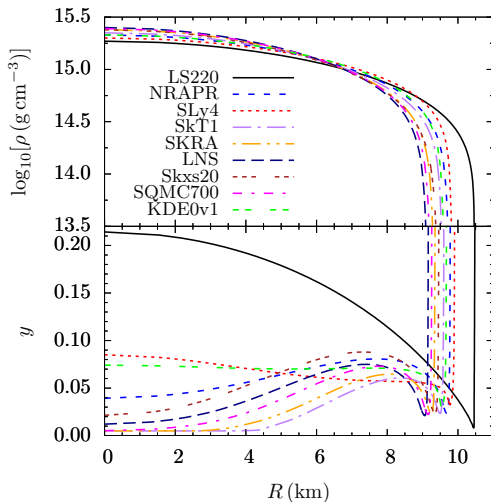
NS mass-radius relationship



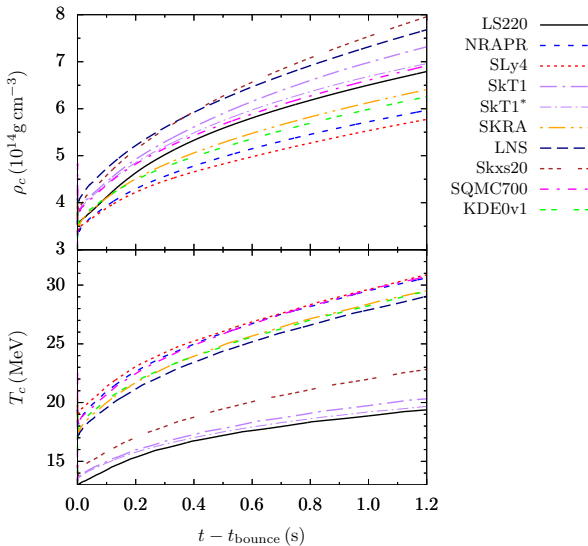
NS Structure

 $1.4M_{\odot}$ 

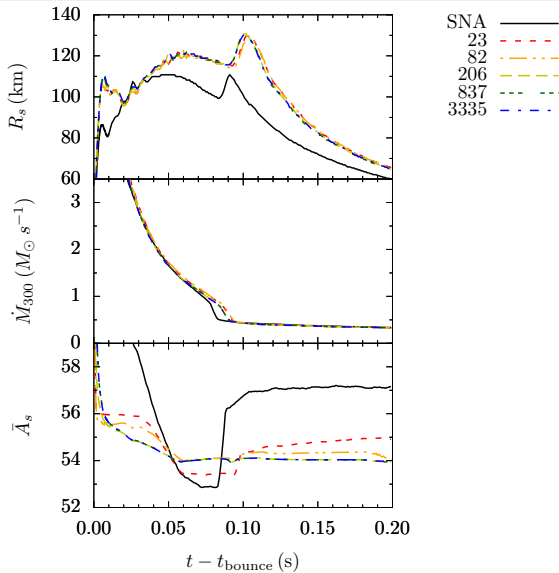
NS Structure

 M_{\max} 

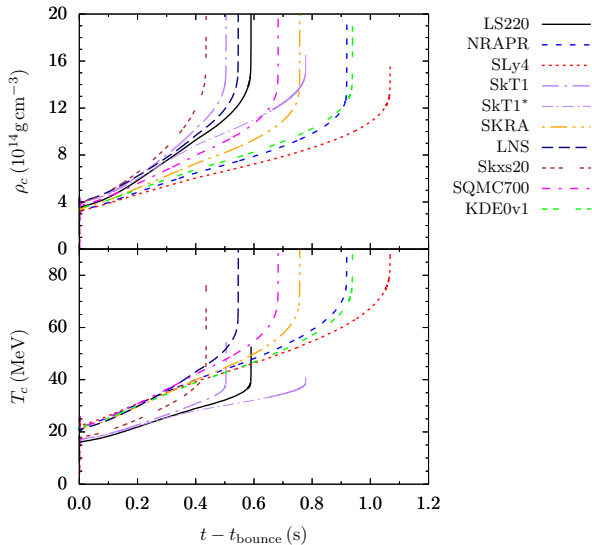
Spherically symmetric collapse of a $15M_{\odot}$ star



Spherically symmetric collapse of a $15M_{\odot}$ star



Spherically symmetric collapse of a $40M_{\odot}$ star



Summary

- Generalize L&S formalism to obtain hot dense EOS for most Skyrme parametrizations.
- Improved calculation of surface properties.
- Added smooth ad-hoc transition from SNA to NSE EOS.
- Extended formalism to allow stiffening of EOS for $n \gtrsim 3n_0$.
- Code converges for large region of parameter space.
 - Temperatures $10^{-4} \text{ MeV} \lesssim T \lesssim 10^{2.5} \text{ MeV}$;
 - Proton fractions $10^{-3} \lesssim y \lesssim 0.70$;
 - Densities $10^{-13} \text{ fm}^{-3} \lesssim n \lesssim 10 \text{ fm}^{-3}$.
- Successfully generated many new EOS tables to study CCSNe, NS mergers, ...

Future

Near future:

- publish results and make code open source;
- study EOS effects on CCSN and NS mergers \Rightarrow EOS may affect neutrino and GW emissions;
- perform 2D and 3D simulations;
- add an improvement treatment of neutron skins to the EOS;
- add reaction network treatment for low temperatures/densities;
- ...